



Beta critical of Schrödinger operator with delta potential in one and two dimension

Rajan Puri

Article

Department of Mathematics and Statistics, Wake Forest University, Winston Salem, NC 27109, USA.; purir@wfu.edu

Academic Editor: Choonkil Park Received: 23 July 2021; Accepted: 22 April 2022; Published: 21 June 2022.

Abstract: We consider 1D and 2D Schrödinger equation with delta potential on the positive half-axis with Dirichlet, Neumann, and Robin type boundary conditions. We presented and estimated the exact values of the beta critical.

Keywords: Schrödinger operator; Delta potential; Critical value; Beta critical; Negative eigenvalues.

MSC: 35J25; 35P15; 47A10; 47D07.

1. Introduction and Main Results

his paper deals with 1D and 2D exactly solvable models for the Schrödinger operators. We study the existence of negative eigenvalues lying under the continuous spectrum of the Schrödinger operator with δ potential. The main goal is to find the threshold value of the coupling constant for which such eigenvalues exist. In one dimensional case, our result's are similar to the consequences of the more general results concerning the Schrödinger operators with a finite number of δ -interactions

$$H = -\frac{d^2}{dx^2} + \sum_{k=1}^n \alpha_k \delta(x - x_k) \tag{1}$$

on the line [1]. In [1], Albeverio and Nizhnik have provided the effective algorithm for determining the number of negative eigenvalues of such operators in terms of the intensities α_k and the distances $d_k = x_{k+1} - x_k$ between the interactions. In this paper, we consider the eigenvalue problem

$$-\frac{d^2y}{dx^2} - \beta\delta(x-a)y(x) = \lambda y(x),$$
(2)

on the positive half-axis with Dirichlet, Neumann, and Robin type boundary conditions. For all solvable models under consideration, there exists a critical value β_{cr} such that the operator possesses the negative eigenvalues if $\beta > \beta_{cr}$ and has no eigenvalues if $\beta > \beta_{cr}$. The goal of this paper is to present and estimate the exact values of the beta critical for Eq. (2) considering appropriate boundary condition. We define beta critical as a critical value of the coupling constant denoted by β_{cr} , the value of β such that Eq. (2) does not have negative eigenvalues for $\beta < \beta_{cr}$ and has them if $\beta > \beta_{cr}$. The delta potential allows solutions for both the bound states $\lambda < 0$ and scattering states $\lambda > 0$. There has been a considerable interest in the study of coupling constant and the problem was investigated by several researchers like S.Albeverio and others in [2], P.Exner and K.Pankrashkin in [3], Barry Simon in [4], Martin Klaus in [5], and Cranston, Koralov, Molchanov and Vainberg in [6]. It is known that the spectrum of $-\Delta - \beta V(x)$ consists of the absolutely continuous part $[0, \infty)$ and at most a finite number of of negative eigenvalues

$$\sigma(-\Delta - \beta V(x)) = \{\lambda_j\} \cup [0, \infty), \ 0 \le j \le N, \quad \lambda_j \le 0.$$
(3)

The classification of the spectrum into discrete and continuous parts usually corresponds to a classification of the dynamics into localized (bound) states and locally decaying states when time increases (scattering) respectively. The lower bound, 0, of the absolutely continuous spectrum is called the ionization threshold. This follows from the fact that the particle is no longer localized, but moves freely when $\lambda > 0$. This classification is related to the space-time behaviour of solutions of the corresponding Schrödinger equation.

It is known that [7] $\beta_{cr} > 0$ in the case of Dirichlet boundary condition and $\beta_{cr} = 0$ in the case of the Neumann boundary condition in the dimension 1 and 2. It was shown that the choice $\beta_{cr} > 0$ or $\beta_{cr} = 0$ depends on whether the truncated resolvent is bounded or goes to infinity when $\lambda \to 0^-$. In fact, β_{cr} was expressed through truncated resolvent operator and depend on the boundary condition and dimension [7]. The main results of this paper in 1D is stated as follows.

Theorem 1. Consider one dimensional eigenvalue problem

$$-\frac{d^2y}{dx^2} - \beta\delta(x-a)y(x) = \lambda y(x), \ y(a) = 1, \ \lambda = -k^2 < 0,$$
(4)

then the $\beta_{cr} = \frac{1}{a} > 0$ in the case of Dirichlet Boundary condition and $\beta_{cr} = 0$ in the case of Neumann and Robin boundary condition.

Proof. The solution of the problem (4) is given by

$$y(x) = Pe^{-kx} + Qe^{kx}, (5)$$

where *P* and *Q* are constants and $\lambda = -k^2 < 0$. One can determine the value of constant P and Q by using the Dirichlet condition y(0) = 0 with y(a) = 1. Then we split the solution of the problem (4) in to two different regions.

$$\begin{cases} y_1(x) = \frac{\sinh kx}{\sinh ka}, & \text{if } 0 \le x \le a\\ y_2(x) = e^{k(a-x)}, & \text{if } a \le x. \end{cases}$$

Integrate the Eq. (4) with respect to *x* over a small interval $\Delta \epsilon$ at the point *x* = *a*,

$$\int_{a-\epsilon}^{a+\epsilon} \left(-\frac{d^2}{dx^2} - \beta \delta(x-a)y \right) dx = \int_{a-\epsilon}^{a+\epsilon} (\lambda y) \, dx. \tag{6}$$

The integral of the second derivative is just the first derivative function and the integral over the function in the right hand side goes to zero. This yields,

$$-y'\Big|_{a-\epsilon}^{a+\epsilon} - \beta y(a) = 0.$$
⁽⁷⁾

When $\epsilon \to 0$, we get

This implies, $\frac{k}{\beta} = \frac{1}{1 + \coth ka}$ and $\frac{ka}{\beta a} = \frac{1}{1 + \coth ka}$. Let ka = A, $\beta a = B$, then $e^{-2A} = 1 - \frac{2A}{B}$. Again let 2A = z, then we have,

 $k + k \coth ka = \beta.$

$$e^{-z} = 1 - \frac{z}{B}.$$
 (8)

From Eq. (8),

and ultimately,
$$\frac{\beta^2}{4} \ge k^2$$
. This leads to

$$\lambda = -k^2 \ge \frac{-\beta^2}{4} \tag{9}$$

Hence, if $\beta = 0$ then there is no possibility of having negative eigenvalues. Therefore, β_{cr} must be greater than zero to produce negative eigenvalues. We notice that we will not have a solution of the Eq. (8) if $\frac{1}{B} \ge 1$. That means, there is no negative eigenvalues when $\frac{1}{a} \ge \beta$. However, for $\frac{1}{B} \le 1$, there is a solution of the Eq. (8). That means, when $\frac{1}{a} \leq \beta$, we will have the negative eigenvalues. As we defined β_{cr} , the value of β such that Eq. (4) does not have negative eigenvalues for $\beta < \beta_{cr}$ and has them if $\beta > \beta_{cr}$, we conclude that $\beta_{cr} = \frac{1}{a}$ for the Eq. (4) with Dirichlet boundary condition.

If we consider the delta potential located at $x = a_n$ then the $\beta_{cr} \rightarrow \infty$ as $a_n \rightarrow 0$. That means, if we approach the potential towards to the boundary, i.e., $a_n \rightarrow 0$ then $\beta_{cr} \rightarrow \infty$.

$$1-\frac{z}{B} \ge 0.$$

 $\frac{\beta}{2} \ge k$

We consider the Neumann boundary condition then the Eq. (2) becomes

$$-\frac{d^2y}{dx^2} - \beta\delta(x-a)y(x) = \lambda y(x), \ y'(0) = 0, \ y(a) = 1, \ \lambda = -k^2 < 0, k > 0.$$
(10)

Similar to Dirichlet problem, we divide the solution of Neumann problem in to two different regions:

$$\begin{cases} y_1(x) = \frac{\cosh kx}{\cosh ka}, & \text{if } 0 \le x \le a\\ y_2(x) = e^{k(a-x)}, & \text{if } a \le x. \end{cases}$$

Integrate the Eq. (10) with respect to *x* over a small interval and take $\epsilon \rightarrow 0$ yields,

$$k + k \tanh ka = \beta$$
.

After simplification, we get

$$\frac{k}{\beta} = \frac{1}{1 + \tanh ka}$$

and then we multiply and divide by *a* for the right hand side,

$$\frac{ka}{\beta a} = \frac{1}{1 + \tanh ka}$$

Let ka = A, $\beta a = B$, then

$$e^{-2A} = \frac{2A}{B} - 1.$$

Following the same analysis as Dirichlet case, we get $\frac{\beta^2}{4} \le k^2$. This implies that $\lambda = -k^2 \le \frac{-\beta^2}{4}$ and concludes that if $\beta = 0$ then there is still a possibility of having a negative eigenvalues. Hence, $\beta_{cr} = 0$.

Now we consider the Robin boundary condition. The Eq. (2) with Robin boundary condition is given by

$$-\frac{d^2y}{dx^2} - \beta\delta(x-a)y(x) = \lambda y(x), \ \frac{dy}{dx} + y(x)\Big|_{x=0} = 0, \ y(a) = 1, \ \lambda = -k^2 < 0.$$
(11)

As above, we divide the solution of this problem in to two different regions:

$$\begin{cases} y_1(x) = \frac{k \cosh kx - \sinh kx}{k \cosh ka - \sinh ka}, & \text{if } 0 \le x \le a \\ y_2(x) = e^{k(a-x)}, & \text{if } a \le x. \end{cases}$$

We integrate the Eq. (11) with respect to *x* over a small interval $\Delta \epsilon$,

$$\int_{a-\epsilon}^{a+\epsilon} (-y'' - \beta \delta(x-a)y) dx = \int_{a-\epsilon}^{a+\epsilon} (\lambda y) dx.$$

After integration and taking $\epsilon \rightarrow 0$, we get,

$$k + k\left(\frac{k - \coth ka}{k \coth ka - 1}\right) = \beta.$$

After simplification,

$$\frac{k}{\beta} = \frac{1}{1 + \left(\frac{k - \coth ka}{k \coth ka - 1}\right)}$$

and

$$\frac{ka}{\beta a} = \frac{1}{1 + \left(\frac{k - \coth ka}{k \coth ka - 1}\right)}.$$

Let ka = A, $\beta a = B$, then

$$e^{-2A} = \frac{2A^2 - 2Aa - AB + Ba}{AB + Ba}$$

Observe that $\frac{2A^2-2Aa-AB+Ba}{AB+Ba}$ must be ≥ 0 . After solving this inequality, we get

$$\frac{\beta}{2} \leq k.$$

Now, $\lambda = -k^2 \leq \frac{-\beta^2}{4}$. Hence, if $\beta = 0$ then there is a possibility of having negative eigenvalues so β_{cr} must be zero. \Box

Two Dimensional eigenvalue problem

The rotational invariance suggests that the two dimensional Laplacian should take a particularly simple form in polar coordinates. We use polar coordinates (r, θ) and look for solutions depending only on r. For d = 2, we do not consider the Neumann boundary condition since β_{cr} is always zero in this case.

Theorem 2. Consider two dimensional Dirichlet problem with delta potential on the circle,

$$-\Delta y(x) - \beta \delta_{1+a} y(x) = \lambda y(x), \ y(1) = 0, \ y(1+a) = 1, \ \lambda = -k^2 < 0, k > 0,$$
(12)

then the $\beta_{cr} > 0$ and $\beta \in (0, 1/2)$.

Proof. The Eq. (2) takes the following form for d = 2 with delta potential on the circle,

$$y'' + \frac{y'}{r} - \beta \delta_{1+a} y(r) = \lambda y(r), y(1) = 0, y(1+a) = 1, \ \lambda = -k^2 < 0, k > 0.$$
⁽¹³⁾

We divide the solution of Eq. (13) into two different regions: region (I) with $1 \le r < 1 + a$ and region (II) with 1 + a < r,

$$\begin{cases} y_1(r) = \frac{Y_0(k)J_0(kr) - J_0(k)Y_0(kr)}{Y_0(k)J_0(k(1+a)) - J_0(k)Y_0(k(1+a))}, & \text{if } 1 \le r \le 1+a\\ y_2(r) = \frac{K_0(kr)}{K_0(k(1+a))}, & \text{if } 1+a \le r. \end{cases}$$

Using the same argument as in the one dimensional problem, we get

$$-y'\Big|_{1+a-\epsilon}^{1+a+\epsilon} -\beta y(1+a) = 0.$$

After simplification, we have

$$-(y_2'|_{1+a+\epsilon}-y_1'|_{1+a-\epsilon})=\beta.$$

When $\epsilon \rightarrow 0$, we get

$$\frac{-Y_0(k)J_1(k(1+a)+J_0(k)Y_1(k(1+a)))}{Y_0(k)J_0(k(1+a))-J_0(k)Y_0(k(1+a))} - \frac{-K_1(k(1+a))}{K_0(k(1+a))} = \frac{\beta}{k},$$
(14)

where Y_0 and Y_1 are Bessel function of second kind and and J_0 and J_1 are Bessel function of first kind. Similarly, K_0 and K_1 are a modified Bessel function of second kind. We define

$$g(k,a) = \frac{g_1(k,a)}{g_2(k,a)} = \frac{-Y_0(k)J_1(k(1+a) + J_0(k)Y_1(k(1+a)))}{Y_0(k)J_0(k(1+a)) - J_0(k)Y_0(k(1+a))}.$$

We notice that g(k, a) = -1 for all the values of *a* as shown in the Figures 1–4.

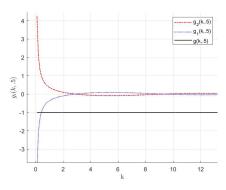


Figure 1. graph of the function g(k,5).

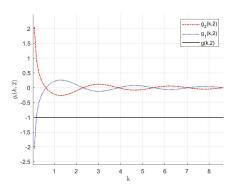


Figure 2. Graph of the function g(k,2).

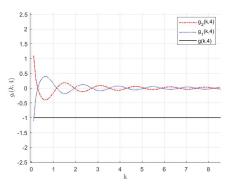


Figure 3. Graph of the function g(k,4).

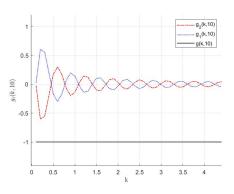


Figure 4. Graph of the function g(k,10).

From Eq. (14), we get

$$-1 + \frac{K_1(k(1+a))}{K_0(k(1+a))} = \frac{\beta}{k}.$$
(15)

We will use the following fact from [8] to prove that $\beta_{cr} > 0$.

Lemma 1 ([8]). Let $p, q \ge 0$. Then the double inequalities

$$1 + \frac{1}{2(x+p)} < \frac{K_1(x)}{K_0(x)} < 1 + \frac{1}{2(x+q)}$$
(16)

hold for all x > 0 if only if $p \ge 1/4$ and q = 0.

Now, from Eqs (15) and (16) we get,

$$\frac{1}{2(x+p)} < \frac{K_1(x)}{K_0(x)} - 1 < \frac{1}{2(x+q)}.$$

When x = k(a + 1) > 0,

$$\frac{1}{2(k(a+1)+p)} < \frac{K_1(k(a+1))}{K_0(k(a+1))} - 1 < \frac{1}{2(k(a+1)+q)}.$$
$$\frac{1}{2(k(a+1)+p)} < \frac{\beta}{k} < \frac{1}{2(k(a+1)+q)}.$$
$$\frac{1}{2((a+1)+\frac{p}{k})} < \beta < \frac{1}{2((a+1)+\frac{q}{k})}.$$

Since a > 0, $p \ge \frac{1}{4}$, k > 0, which tells us that $\beta > 0$ and hence $\beta_{cr} > 0$ and $\beta \in (0, \frac{1}{2})$. \Box

Conflicts of Interest: "The author declares no conflict of interest."

Data Availability: All data required for this research is included within this paper.

Funding Information: No funding is available for this research.

References

- Albeverio, S., & Nizhnik, L. (2003). On the number of negative eigenvalues of a one-dimensional Schrödinger operator with point interactions. *Letters in Mathematical Physics*, 65(1), 27-35.
- [2] Albeverio, S., Gesztesy, F., Hoegh-Krohn, R., & Holden, H. (2012). Solvable Models in Quantum Mechanics. Springer Science & Business Media.
- [3] Exner, P., & Pankrashkin, K. (2014). Strong coupling asymptotics for a singular Schrödinger operator with an interaction supported by an open arc. *Communications in Partial Differential Equations*, 39(2), 193-212.
- [4] Simon, B. (1976). The bound state of weakly coupled Schrödinger operators in one and two dimensions. *Annals of Physics*, 97(2), 279-288.
- [5] Klaus, M. (1977). On the bound state of Schrödinger operators in one dimension. Annals of Physics, 108(2), 288-300.
- [6] Cranston, M., Koralov, L., Molchanov, S., & Vainberg, B. (2009). Continuous model for homopolymers. *Journal of Functional Analysis*, 256(8), 2656-2696.
- [7] Puri, R., & Vainberg, B. (2022). On the critical value of the coupling constant in exterior elliptic problems. *Applicable Analysis*, 101(1), 108-117.
- [8] Yang, Z. H., & Chu, Y. M. (2017). On approximating the modified Bessel function of the second kind. *Journal of Inequalities and Applications*, 2017, Article No 47. https://doi.org/10.1186/s13660-017-1317-z.



© 2022 by the authors; licensee PSRP, Lahore, Pakistan. This article is an open access article distributed under the terms and conditions of the Creative Commons Attribution (CC-BY) license (http://creativecommons.org/licenses/by/4.0/).